

Fourier Series

Chapter Intended Learning Outcomes:

- (i) Represent continuous-time periodic signals using Fourier series
- (ii) Understand the properties of Fourier series
- (iii) Understand the relationship between Fourier series and linear time-invariant system

Periodic Signal Representation in Frequency Domain

Fourier series can be considered as the frequency domain representation of a continuous-time periodic signal.

Recall (2.6) that $x(t)$ is said to be periodic if there exists $T_p > 0$ such that

$$x(t) = x(t + T_p), \quad t \in (-\infty, \infty) \quad (4.1)$$

The smallest T_p for which (4.1) holds is called the fundamental period.

Using (2.26), the fundamental frequency is related to T_p as:

$$\Omega_0 = \frac{2\pi}{T_p} \quad (4.2)$$

According to Fourier series, $x(t)$ is represented as:

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\Omega_0 t}, \quad t \in (-\infty, \infty) \quad (4.3)$$

where

$$a_k = \frac{1}{T_p} \int_{-T_p/2}^{T_p/2} x(t) e^{-jk\Omega_0 t} dt, \quad k = \dots - 1, 0, 1, 2, \dots \quad (4.4)$$

are called **Fourier series coefficients**. Note that the integration can be done for any period, e.g., $(0, T_p)$, $(-T_p, 0)$.

That is, every **periodic** signal can be expressed as a sum of **harmonically related complex sinusoids** with frequencies $\dots - \Omega_0, 0, \Omega_0, 2\Omega_0, 3\Omega_0, \dots$, where the fundamental frequency Ω_0 is called the **first harmonic**, $2\Omega_0$ is called the **second harmonic**, and so on.

This means that $x(t)$ only contains frequencies $\dots - \Omega_0, 0, \Omega_0, 2\Omega_0, \dots$ with 0 being the DC component.

Note that the sinusoids are **complex-valued** with both **positive** and **negative** frequencies.

Note also that a_k is generally **complex** and we can also use magnitude and phase for its representation:

$$|a_k| = \sqrt{(\Re\{a_k\})^2 + (\Im\{a_k\})^2} \quad (4.5)$$

and

$$\angle(a_k) = \tan^{-1} \left(\frac{\Im\{a_k\}}{\Re\{a_k\}} \right) \quad (4.6)$$

From (4.3), $\{a_k\}$ can be used to represent $x(t)$.

Example 4.1

Find the Fourier series coefficients for $x(t) = \cos(10\pi t) + \cos(20\pi t)$.

It is clear that the fundamental frequency of $x(t)$ is $\Omega_0 = 10\pi$. According to (4.2), the fundamental period is thus equal to $T_p = 2\pi/\Omega_0 = 1/5$, which is validated as follows:

$$\begin{aligned}x\left(t + \frac{1}{5}\right) &= \cos\left(10\pi\left(t + \frac{1}{5}\right)\right) + \cos\left(20\pi\left(t + \frac{1}{5}\right)\right) \\ &= \cos(10\pi t + 2\pi) + \cos(20\pi t + 4\pi) \\ &= \cos(10\pi t) + \cos(20\pi t)\end{aligned}$$

With the use of Euler formula in (2.29):

$$\cos(u) = \frac{e^{ju} + e^{-ju}}{2}$$

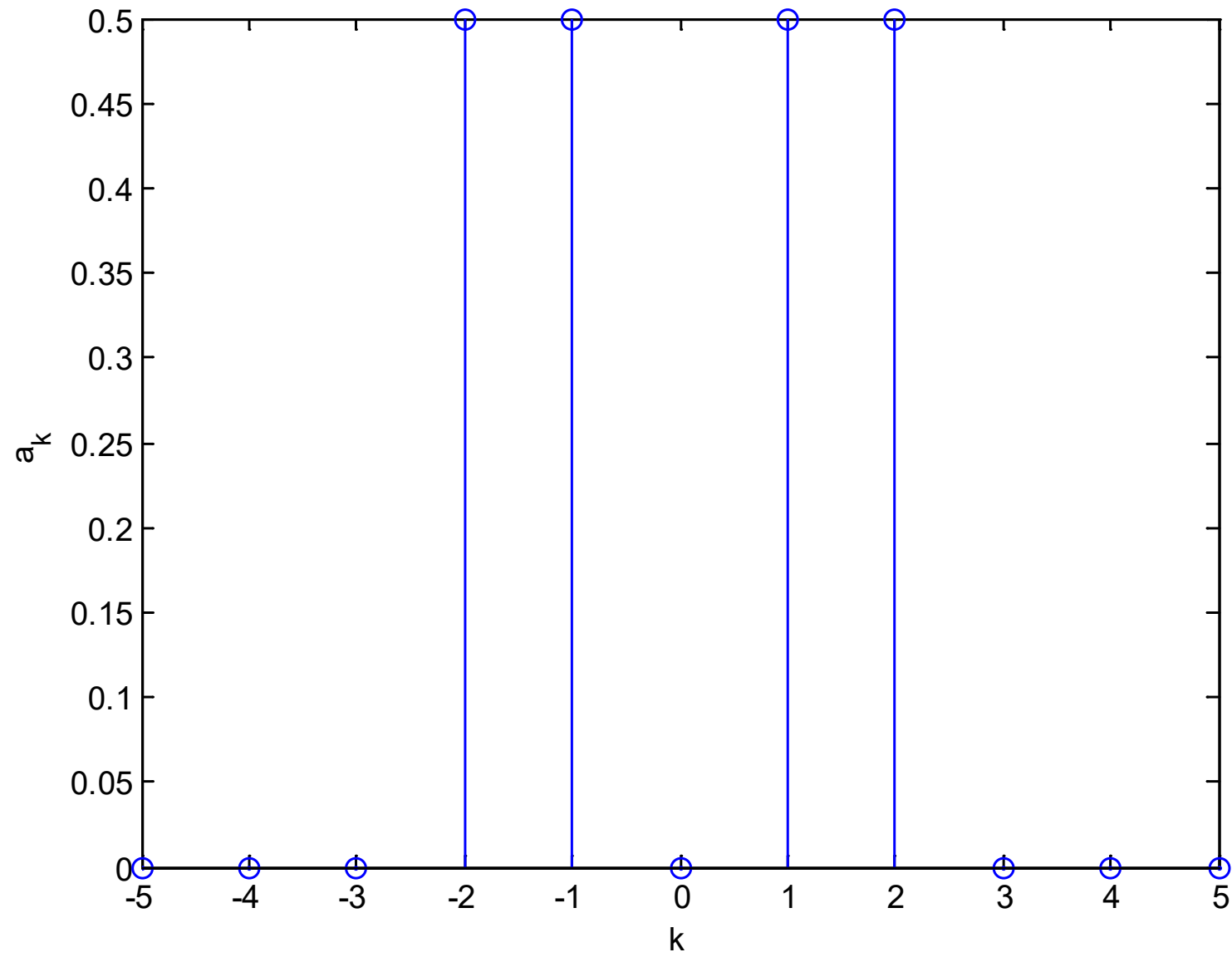
We can express $x(t)$ as:

$$\begin{aligned}x(t) &= \cos(10\pi t) + \cos(20\pi t) \\&= \frac{e^{j10\pi t} + e^{-j10\pi t}}{2} + \frac{e^{j20\pi t} + e^{-j20\pi t}}{2} \\&= \frac{1}{2}e^{-j20\pi t} + \frac{1}{2}e^{-j10\pi t} + \frac{1}{2}e^{j10\pi t} + \frac{1}{2}e^{j20\pi t}\end{aligned}$$

which only contains four frequencies. Comparing with (4.3):

$$a_k = \begin{cases} 0.5, & k = -2 \\ 0.5, & k = -1 \\ 0.5, & k = 1 \\ 0.5, & k = 2 \\ 0, & \text{otherwise} \end{cases}$$

Can we use (4.4)? Why?



Example 4.2

Find the Fourier series coefficients for $x(t) = 1 + \sin(\Omega_0 t) + 2 \cos(\Omega_0 t) + \cos(3\Omega_0 t + \pi/4)$.

With the use of Euler formulas in (2.29)-(2.30), $x(t)$ can be written as:

$$\begin{aligned} x(t) &= 1 + \left(1 + \frac{1}{2j}\right) e^{j\Omega_0 t} + \left(1 - \frac{1}{2j}\right) e^{-j\Omega_0 t} + \frac{1}{2} e^{j\pi/4} e^{3j\Omega_0 t} + \frac{1}{2} e^{-j\pi/4} e^{-3j\Omega_0 t} \\ &= \frac{\sqrt{2}}{4} (1 - j) e^{-3j\Omega_0 t} + \left(1 + j\frac{1}{2}\right) e^{-j\Omega_0 t} + 1 + \left(1 - j\frac{1}{2}\right) e^{j\Omega_0 t} \\ &\quad + \frac{\sqrt{2}}{4} (1 + j) e^{3j\Omega_0 t} \end{aligned}$$

Again, comparing with (4.3) yields:

$$a_k = \begin{cases} \frac{\sqrt{2}}{4}(1 - j), & k = -3 \\ 1 + \frac{j}{2}, & k = -1 \\ 1, & k = 0 \\ 1 - \frac{j}{2}, & k = 1 \\ \frac{\sqrt{2}}{4}(1 + j), & k = 3 \\ 0, & \text{otherwise} \end{cases}$$

To plot $\{a_k\}$, we may compute $|a_k|$ and $\angle(a_k)$ for all k , e.g.,

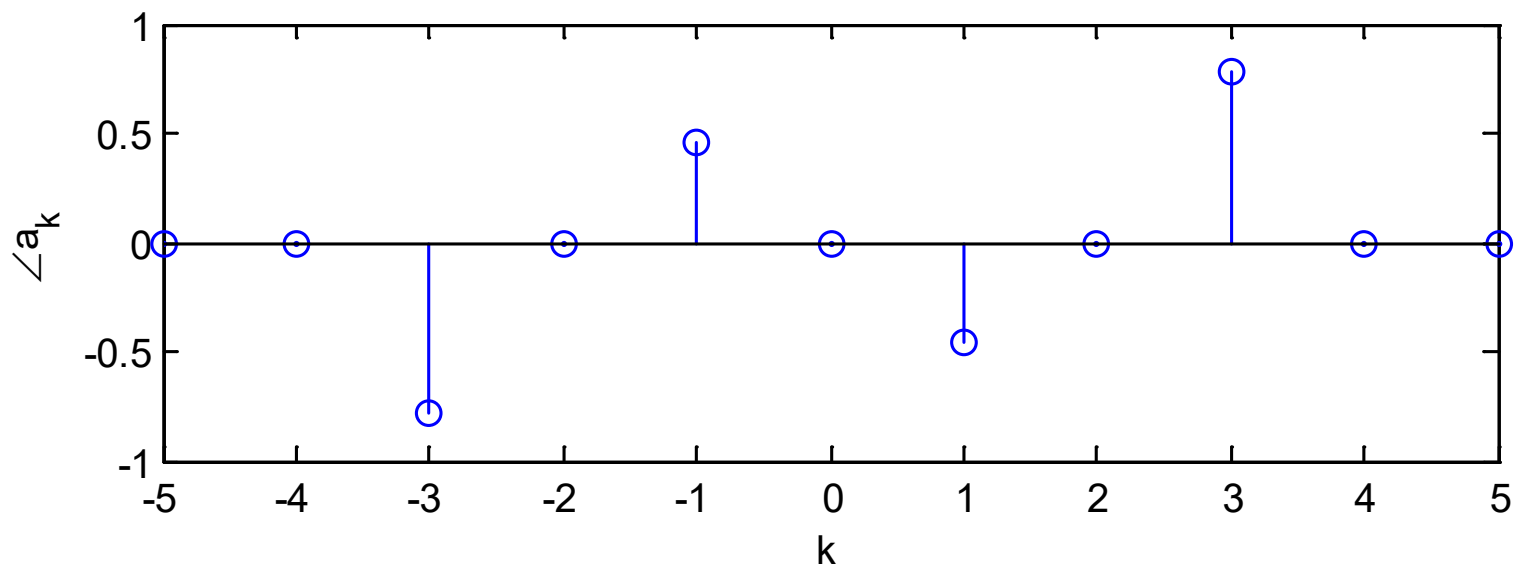
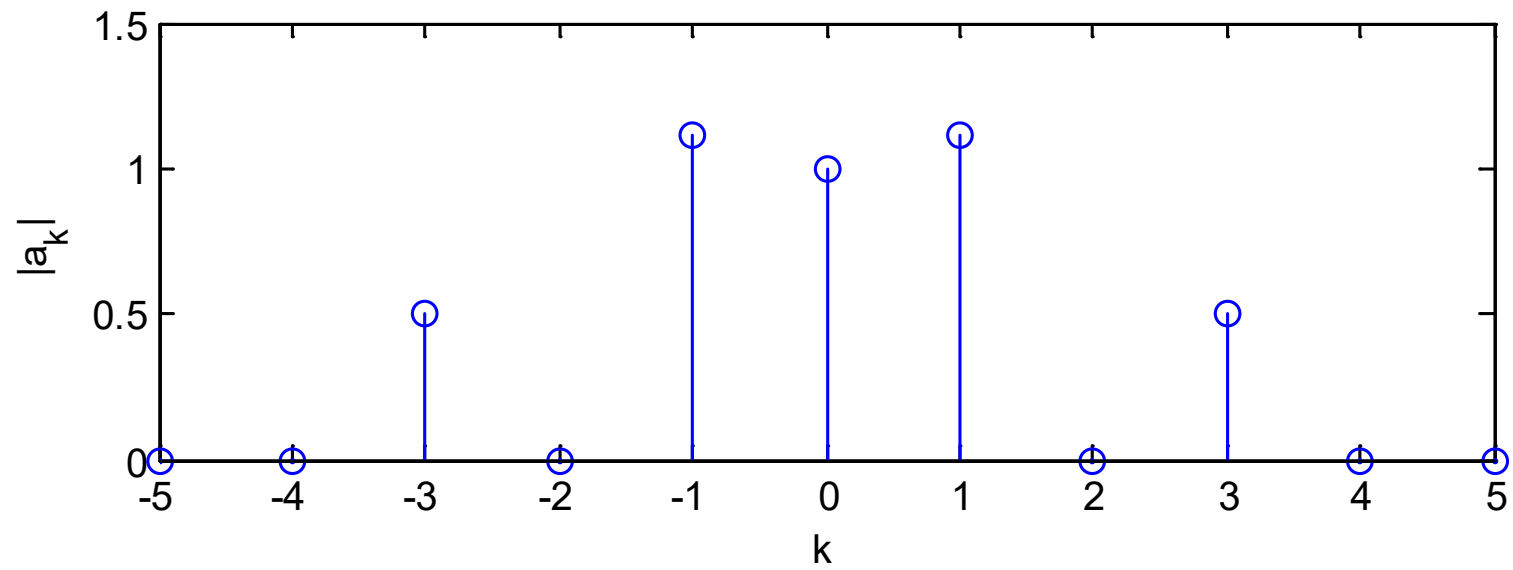
$$|a_{-3}| = \sqrt{\left(\frac{\sqrt{2}}{4}\right)^2 + \left(-\frac{\sqrt{2}}{4}\right)^2} = \frac{1}{2}$$

and

$$\angle(a_{-3}) = \tan^{-1}(-1) = -\frac{\pi}{4}$$

We can also use MATLAB commands `abs` and `angle` to compute the magnitude and phase, respectively. After constructing a vector \mathbf{x} containing $\{a_k\}$, we can plot $|a_k|$ and $\angle(a_k)$ using:

```
subplot(2,1,1)
stem(n,abs(x))
xlabel('k')
ylabel('|a_k|')
subplot(2,1,2)
stem(n,angle(x))
xlabel('k')
ylabel('\angle{a_k}')
```

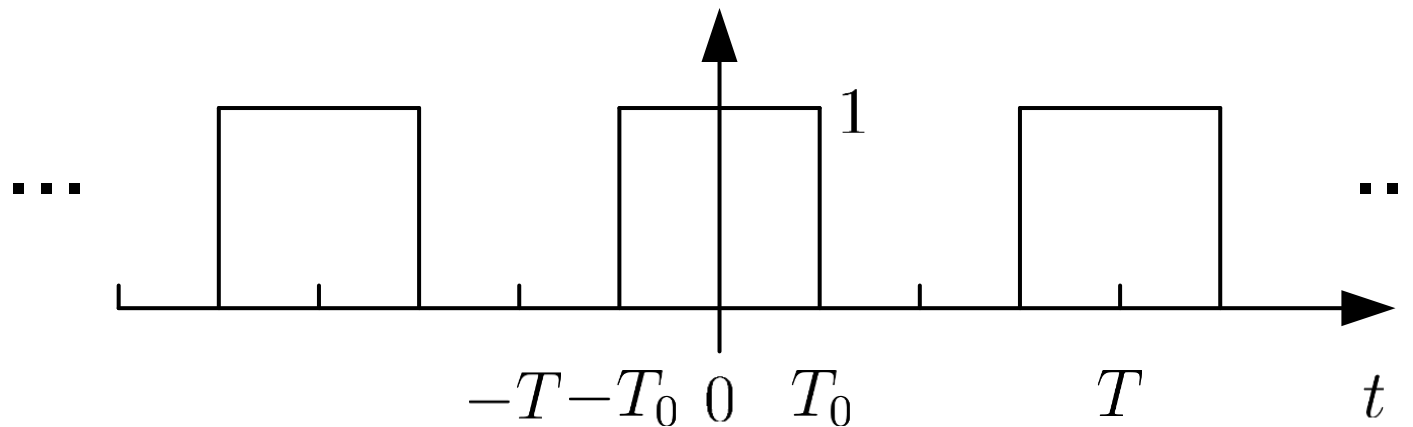


Example 4.3

Find the Fourier series coefficients for $x(t)$, which is a periodic continuous-time signal of fundamental period T and is a pulse with a width of $2T_0$ in each period. Over the specific period from $-T/2$ to $T/2$, $x(t)$ is:

$$x(t) = \begin{cases} 1, & -T_0 < t < T_0 \\ 0, & \text{otherwise} \end{cases}$$

with $T > 2T_0$.



Noting that the fundamental frequency is $\Omega_0 = 2\pi/T$ and using (4.4), we get:

$$a_k = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-jk\Omega_0 t} dt = \frac{1}{T} \int_{-T_0}^{T_0} e^{-jk\Omega_0 t} dt$$

For $k = 0$:

$$a_0 = \frac{1}{T} \int_{-T_0}^{T_0} 1 dt = \frac{2T_0}{T}$$

For $k \neq 0$:

$$a_k = \frac{1}{T} \int_{-T_0}^{T_0} e^{-jk\Omega_0 t} dt = -\frac{1}{jk\Omega_0 T} e^{-jk\Omega_0 t} \Big|_{-T_0}^{T_0} = \frac{\sin(k\Omega_0 T_0)}{k\pi} = \frac{\sin(2\pi k T_0 / T)}{k\pi}$$

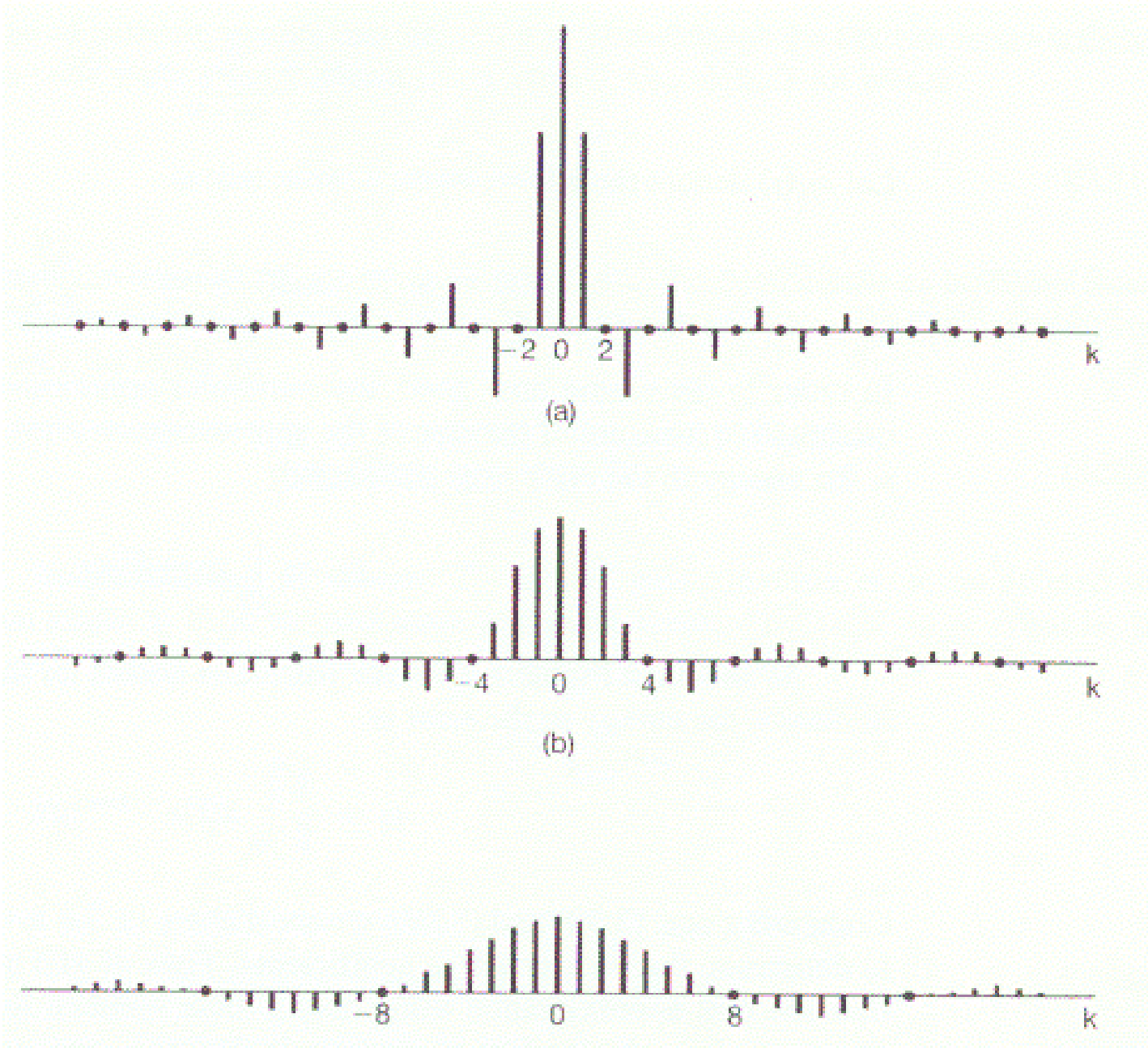
The reason of separating the cases of $k = 0$ and $k \neq 0$ is to facilitate the computation of a_0 , whose value is not straightforwardly obtained from the general expression which involves "0/0".

Nevertheless, using L'Hôpital's rule:

$$\lim_{k \rightarrow 0} \frac{\sin(2\pi k T_0/T)}{k\pi} = \lim_{k \rightarrow 0} \frac{\frac{d \sin(2\pi k T_0/T)}{dk}}{\frac{dk\pi}{dk}} = \lim_{k \rightarrow 0} \frac{2\pi T_0/T \cos((2\pi k T_0/T))}{\pi} = \frac{2T_0}{T}$$

An investigation on the values of $\{a_k\}$ with respect to relative pulse width T_0/T is performed as follows.

We see that when T_0/T decreases, $\{a_k\}$ seem to be stretched.



Example 4.4

Find the Fourier series coefficients for the following continuous-time periodic signal $x(t)$:

$$x(t) = \begin{cases} 1.5, & 0 < t < 1 \\ -1.5, & 1 < t < 2 \end{cases}$$

where the fundamental period is $T_p = 2$ and fundamental frequency is $\Omega_0 = \pi$.

Using (4.4) with the period from $t = -1$ to $t = 1$:

$$\begin{aligned} a_k &= \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-jk\Omega_0 t} dt \\ &= \frac{1}{2} \int_{-1}^0 (-1.5) e^{-jk\pi t} dt + \frac{1}{2} \int_0^1 1.5 e^{-jk\pi t} dt \end{aligned}$$

For $k = 0$:

$$a_k = \frac{1}{2} \int_{-1}^0 (-1.5) dt + \frac{1}{2} \int_0^1 1.5 dt = \frac{1}{2} (-1.5 + 1.5) = 0$$

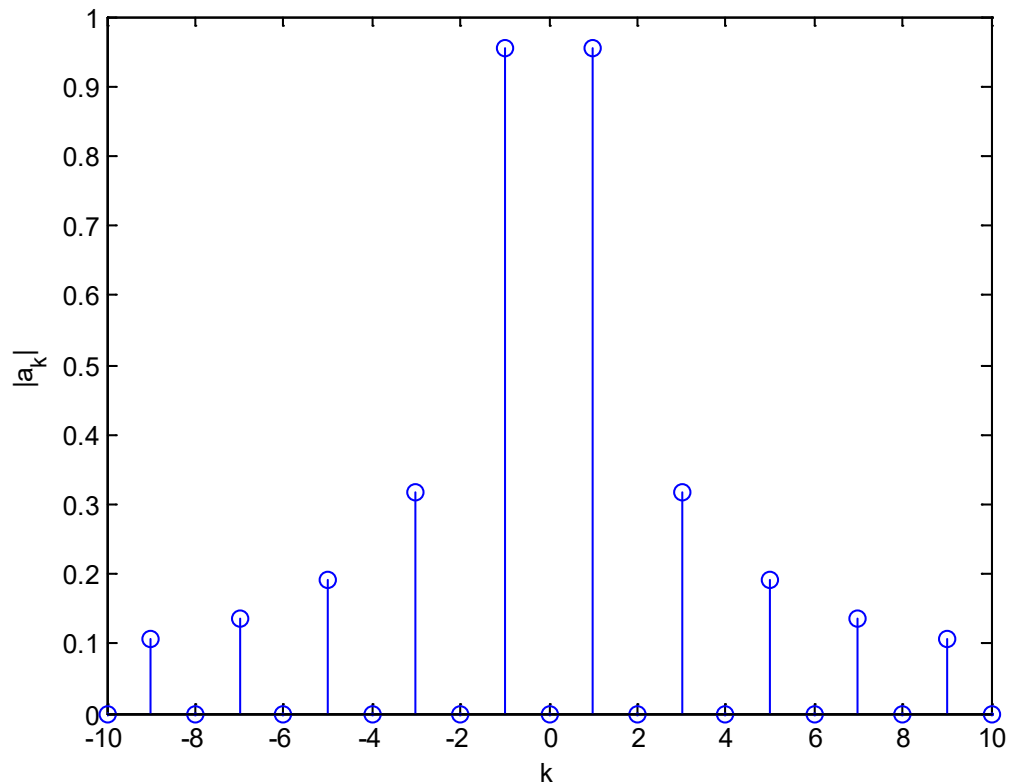
For $k \neq 0$:

$$\begin{aligned} a_k &= \frac{1}{2} \int_{-1}^0 (-1.5) e^{-jk\pi t} dt + \frac{1}{2} \int_0^1 1.5 e^{-jk\pi t} dt \\ &= \frac{3}{4} \left[\int_{-1}^0 -e^{-jk\pi t} dt + \int_0^1 e^{-jk\pi t} dt \right] \\ &= \frac{3}{4} \left[-\frac{1}{-jk\pi} e^{-jk\pi t} \Big|_{-1}^0 + \frac{1}{-jk\pi} -e^{-jk\pi t} \Big|_0^1 \right] \\ &= \frac{3}{4jk\pi} [1 - e^{jk\pi} - e^{-jk\pi} + 1] \\ &= \frac{3}{2jk\pi} [1 - \cos(k\pi)] \end{aligned}$$

MATLAB can be used to validate the answer. First we have:

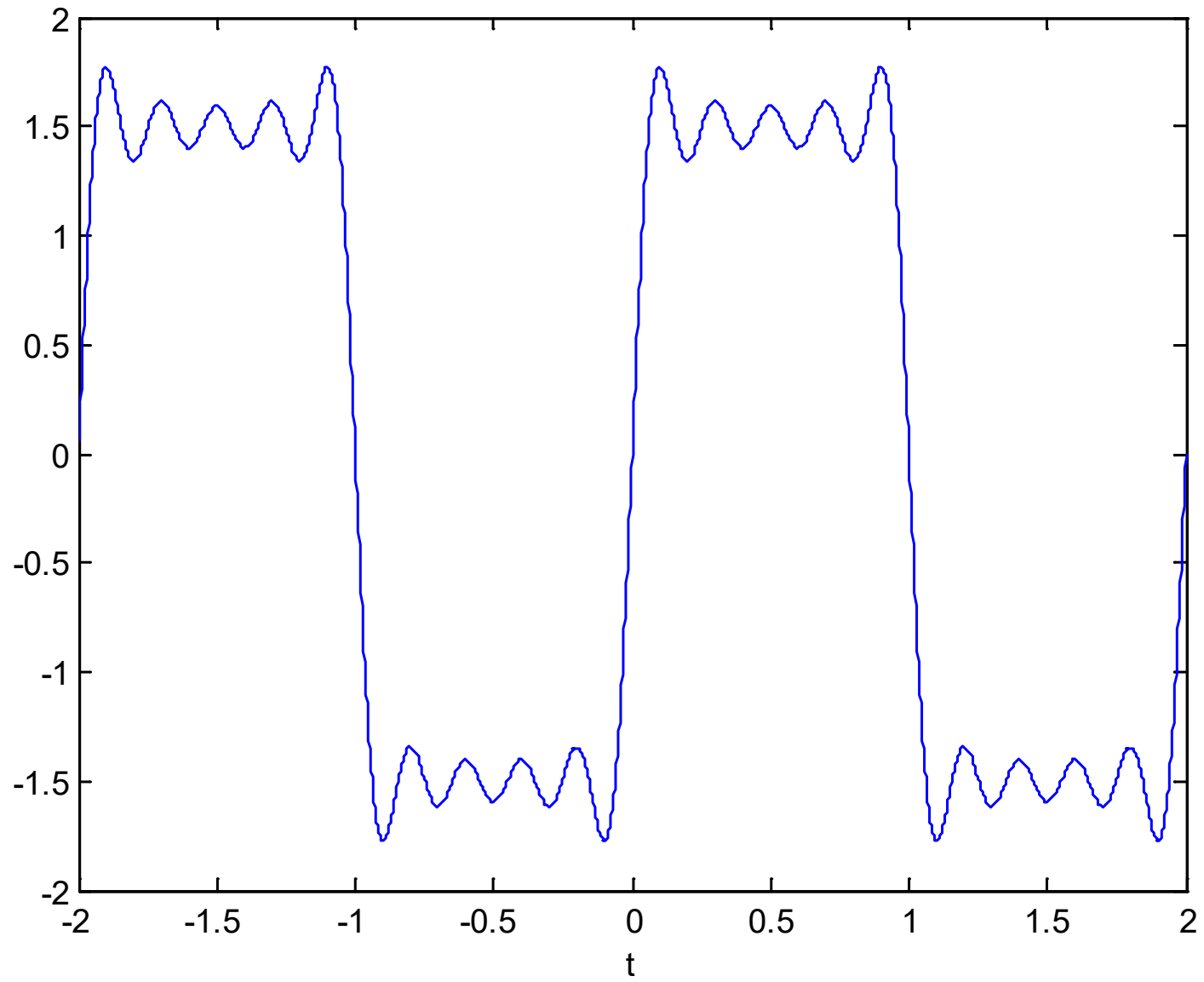
$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\Omega_0 t} \approx \sum_{k=-K}^K a_k e^{jk\Omega_0 t}$$

for sufficiently large K because $|a_k|$ is decreasing with k

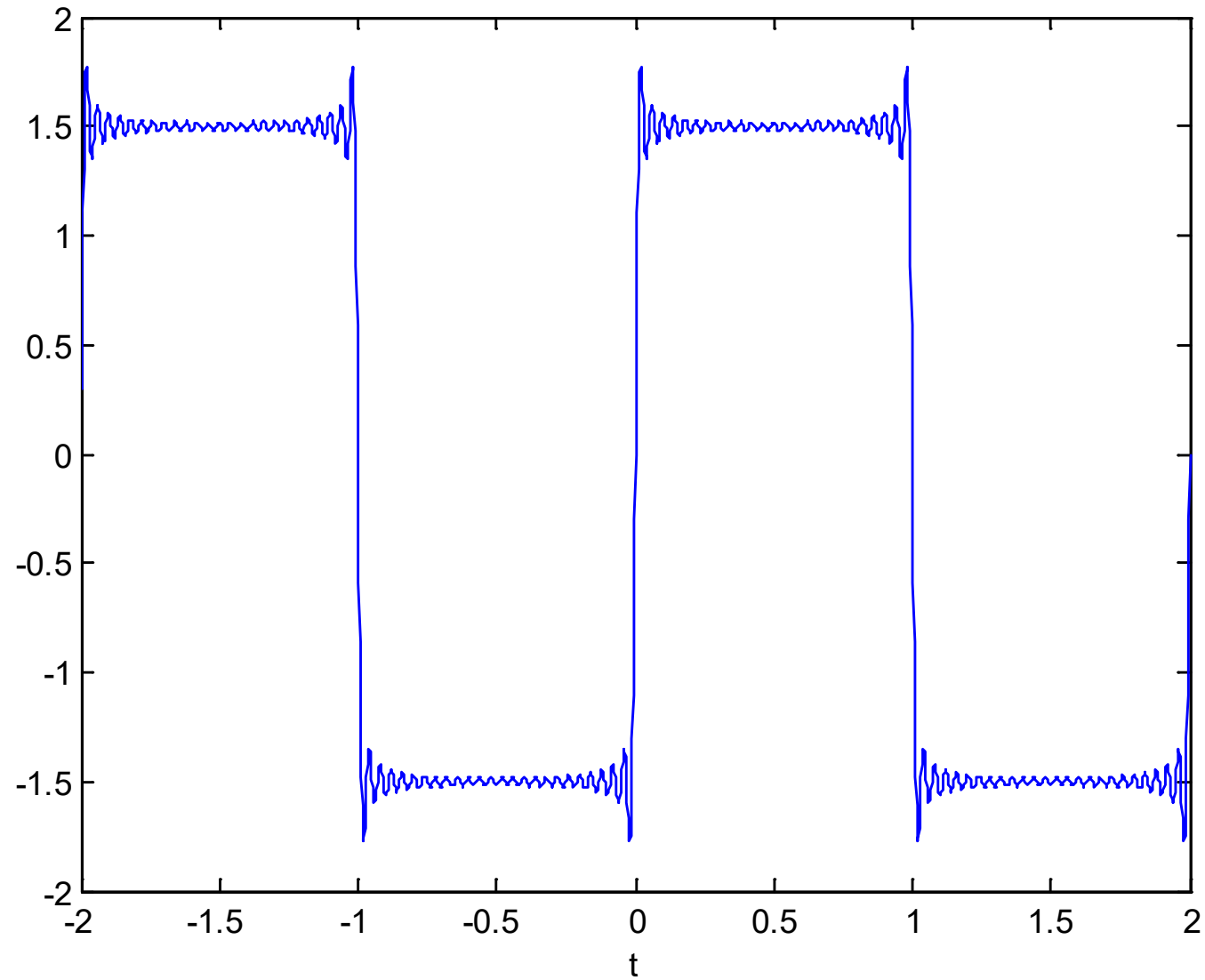


Setting $K = 10$, we may use the following code:

```
K=10;
a_p = 3./(j.*2.*[1:K].*pi).*(1-cos([1:K].*pi)); % +ve a_k
a_n = 3./(j.*2.*[-K:-1].*pi).*(1-cos([-K:-1].*pi)); %-ve a_k
a = [a_n 0 a_p]; %construct vector of a_k
for n=1:2000
    t=(n-1000)/500; %time interval of (-2,2);
        %small sampling interval of 1/500 to approximate x(t);
    e = (exp(j.*[-K:K].*pi.*t)).'; %construct exponential vector
    x(n) = a*e;
end
x=real(x); %remove imaginary parts due to precision error
n=1:2000;
t=(n-1000)./500;
plot(t,x)
xlabel('t')
```



For $K = 50$:



In summary, if $x(t)$ is periodic, it can be represented as a **linear** combination of complex harmonics with amplitudes $\{a_k\}$.

That is, $\{a_k\}$ correspond to the frequency domain representation of $x(t)$ and we may write:

$$x(t) \leftrightarrow X(j\Omega) \quad \text{or} \quad x(t) \leftrightarrow a_k \quad (4.7)$$

where $X(j\Omega)$, a function of frequency Ω , is characterized by $\{a_k\}$.

Both $x(t)$ and $X(j\Omega)$ represent the **same** signal: we observe the former in time domain while the latter in frequency domain.

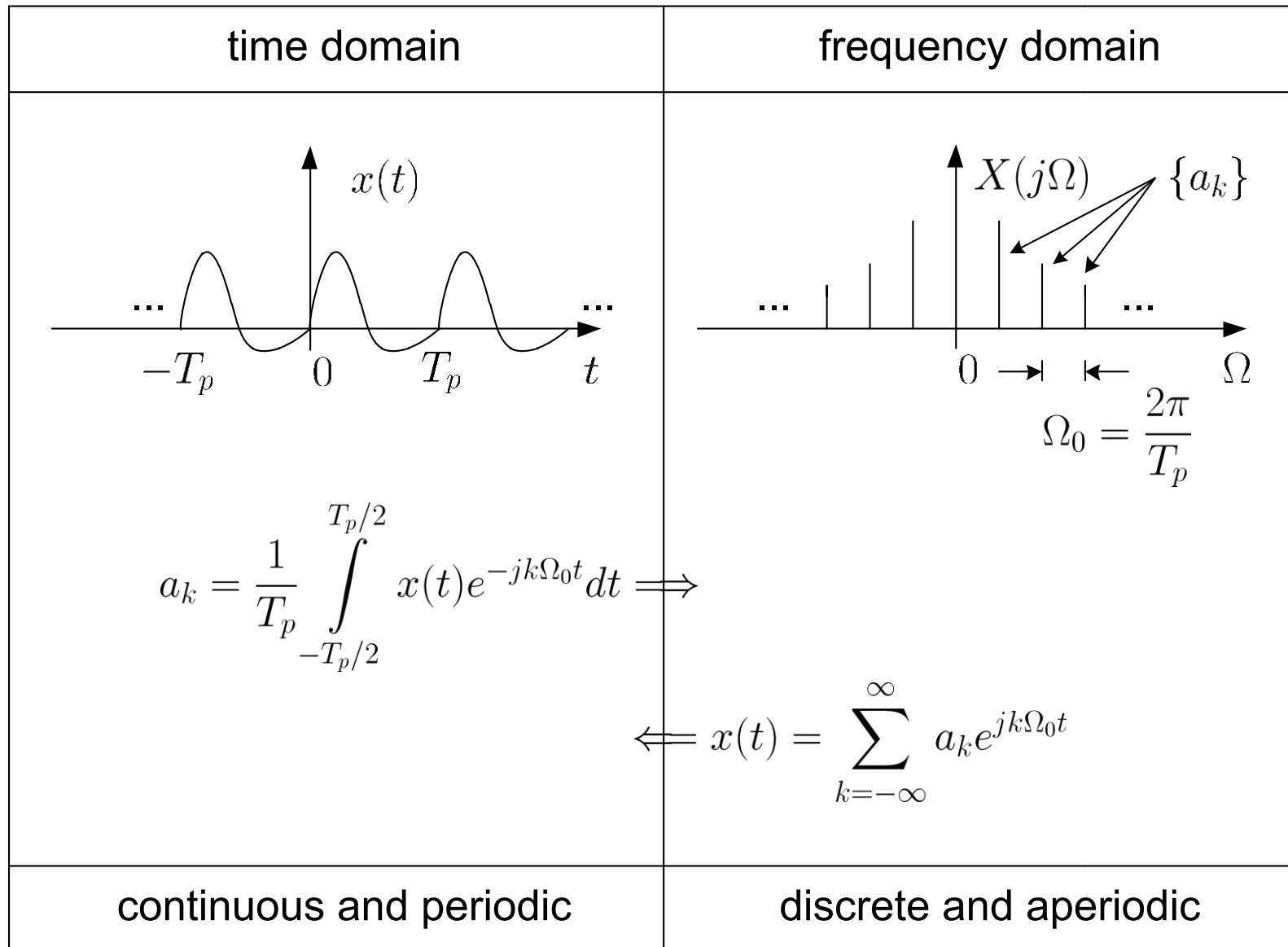


Fig.4.1: Illustration of Fourier series

Properties of Fourier Series

Linearity

Let $x(t) \leftrightarrow a_k$ and $y(t) \leftrightarrow b_k$ be two Fourier series pairs with the same period of T_p . We have:

$$Ax(t) + By(t) \leftrightarrow Aa_k + Bb_k \quad (4.8)$$

This can be proved as follows. As $x(t)$ and $y(t)$ have the same fundamental period of T_p or fundamental frequency Ω_0 , we can write:

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\Omega_0 t}, \quad y(t) = \sum_{k=-\infty}^{\infty} b_k e^{jk\Omega_0 t}$$

Multiplying $x(t)$ and $y(t)$ by A and B , respectively, yields:

$$Ax(t) = A \sum_{k=-\infty}^{\infty} a_k e^{jk\Omega_0 t}, \quad By(t) = B \sum_{k=-\infty}^{\infty} b_k e^{jk\Omega_0 t}$$

Summing $Ax(t)$ and $By(t)$, we get:

$$Ax(t) + By(t) = \sum_{k=-\infty}^{\infty} (Aa_k + Bb_k) e^{jk\Omega_0 t} \leftrightarrow Aa_k + Bb_k$$

Time Shifting

A shift of t_0 in $x(t)$ causes a multiplication of $e^{-jk\Omega_0 t_0}$ in a_k :

$$x(t) \leftrightarrow a_k \Rightarrow x(t - t_0) \leftrightarrow e^{-jk\Omega_0 t_0} a_k = e^{-jk(2\pi)/T_p t_0} a_k \quad (4.9)$$

Time Reversal

$$x(t) \leftrightarrow a_k \Rightarrow x(-t) \leftrightarrow a_{-k} \quad (4.10)$$

(4.9) and (4.10) are proved as follows.

Recall (4.3):

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\Omega_0 t}$$

Substituting t by $t - t_0$, we obtain:

$$x(t - t_0) = \sum_{k=-\infty}^{\infty} a_k e^{jk\Omega_0(t-t_0)} = \sum_{k=-\infty}^{\infty} (e^{-jk\Omega_0 t_0} a_k) e^{jk\Omega_0 t} \leftrightarrow e^{-jk\Omega_0 t_0} a_k$$

Substituting t by $-t$ yields:

$$x(-t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\Omega_0(-t)} = \sum_{l=-\infty}^{\infty} a_{-l} e^{jl\Omega_0 t} = \sum_{k=-\infty}^{\infty} a_{-k} e^{jk\Omega_0 t} \leftrightarrow a_{-k}$$

Time Scaling

For a time-scaled version of $x(t)$, $x(\alpha t)$ where $\alpha > 0$ is a real number, we have:

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\Omega_0 t} \Rightarrow x(\alpha t) = \sum_{k=-\infty}^{\infty} a_k e^{jk(\alpha\Omega_0)t} \quad (4.11)$$

Multiplication

Let $x(t) \leftrightarrow a_k$ and $y(t) \leftrightarrow b_k$ be two Fourier series pairs with the same period of T_p . We have:

$$x(t)y(t) \leftrightarrow \sum_{l=-\infty}^{\infty} a_l b_{k-l} = a_k \otimes b_k \quad (4.12)$$

(4.12) is proved as follows.

Applying (4.3) again, the product of $x(t)$ and $y(t)$ is:

$$\begin{aligned}x(t)y(t) &= \sum_{l=-\infty}^{\infty} a_l e^{jl\Omega_0 t} \sum_{n=-\infty}^{\infty} b_n e^{jn\Omega_0 t} \\&= \sum_{l=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} a_l b_n e^{j(l+n)\Omega_0 t} \\&= \sum_{l=-\infty}^{\infty} \sum_{k=-\infty}^{\infty} a_l b_{l-k} e^{jk\Omega_0 t}, \quad k = l + n \\&= \sum_{k=-\infty}^{\infty} \left(\sum_{l=-\infty}^{\infty} a_l b_{l-k} \right) e^{jk\Omega_0 t} \leftrightarrow \sum_{l=-\infty}^{\infty} a_l b_{k-l}\end{aligned}$$

Conjugation

$$x(t) \leftrightarrow a_k \Rightarrow x^*(t) \leftrightarrow a_{-k}^* \quad (4.13)$$

Parseval's Relation

The Parseval's relation addresses the **power** of $x(t)$:

$$\frac{1}{T_p} \int_{-T_p/2}^{T_p/2} |x(t)|^2 dt = \sum_{k=-\infty}^{\infty} |a_k|^2 \quad (4.14)$$

That is, we can compute the power in either the time domain or frequency domain.

Example 4.5

Prove the Parseval's relation.

Using (4.3), we have:

$$\begin{aligned}\frac{1}{T_p} \int_{-T_p/2}^{T_p/2} |x(t)|^2 dt &= \frac{1}{T_p} \int_{-T_p/2}^{T_p/2} \left(\sum_{m=-\infty}^{\infty} a_m e^{jm\Omega_0 t} \right) \left(\sum_{n=-\infty}^{\infty} a_n e^{jn\Omega_0 t} \right)^* dt \\ &= \frac{1}{T_p} \int_{-T_p/2}^{T_p/2} \left(\sum_{m=-\infty}^{\infty} a_m e^{jm\Omega_0 t} \right) \left(\sum_{n=-\infty}^{\infty} a_n^* e^{-jn\Omega_0 t} \right) dt \\ &= \frac{1}{T_p} \sum_{m=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} a_m a_n^* \int_{-T_p/2}^{T_p/2} e^{j(m-n)\Omega_0 t} dt \\ &= \frac{1}{T_p} \sum_{m=-\infty}^{\infty} |a_m|^2 \int_{-T_p/2}^{T_p/2} dt = \sum_{m=-\infty}^{\infty} |a_m|^2\end{aligned}$$

Linear Time-Invariant System with Periodic Input

Recall for a linear time-invariant (LTI) system, the input-output relationship is characterized by convolution in (3.17):

$$\begin{aligned} y(t) &= x(t) \otimes h(t) \\ &= \int_{-\infty}^{\infty} x(\tau)h(t - \tau)d\tau = \int_{-\infty}^{\infty} h(\tau)x(t - \tau)d\tau \end{aligned} \quad (4.15)$$

If the input to the system with impulse response $h(t)$ is $x(t) = e^{j\Omega_0 t}$, then the output is:

$$\begin{aligned} y(t) &= \int_{-\infty}^{\infty} h(\tau)x(t - \tau)d\tau \\ &= \int_{-\infty}^{\infty} h(\tau)e^{j\Omega_0(t-\tau)}d\tau \\ &= e^{j\Omega_0 t} \int_{-\infty}^{\infty} h(\tau)e^{-j\Omega_0\tau}d\tau \end{aligned} \quad (4.16)$$

Note that $\int_{-\infty}^{\infty} h(\tau)e^{-j\Omega_0\tau}d\tau$ is independent of t but a function of Ω_0 and we may denote it as $H(j\Omega_0)$:

$$y(t) = e^{j\Omega_0 t} H(j\Omega_0) = H(j\Omega_0)x(t) \quad (4.17)$$

If we input a sinusoid through a LTI system, there is **no change in frequency** in the output but amplitude and phase are modified.

Generalizing the result to any periodic signal in (4.3) yields:

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\Omega_0 t} \rightarrow y(t) = \sum_{k=-\infty}^{\infty} a_k H(jk\Omega_0) e^{jk\Omega_0 t} \quad (4.18)$$

where only the Fourier series coefficients are modified.

Note that discrete Fourier series is used to represent discrete periodic signal in (2.7) but it will not be discussed.